

Analysis of Fundamental Photodetection Noises and Evaluation of PIN and APD Photodiodes Performances using an Optical High Debit Transmission Chain Simulated by Optisystem

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ABSTRACT

Regarded as an extremity component for an optical communication chain, the photodetector plays a very significant role to detect and convert the optical ray penetrated by the optical fiber into an electric signal easier to exploit in the extraction of the information. Indeed, the noise of photodetection constitutes a precious phenomenon having several physical sources that affects the delivered signal at the output of the photodetector, focusing fundamentally on thermal noise “*Johnson noise*”, quantum noise “*Shot noise*” and noise of multiplication “*Multiplicative noise*”. This article aims to present a total study about the physical origins of fundamental photodetection noises by analyzing their behavior in *PIN* and *APD* photodiodes through an optical high debit transmission chain simulated by *Optisystem* software. The aim of this research extends also to evaluate and compare between *PIN* and *APD* photodiodes in function of their functional parameters in terms of performance by measurement of signal-to-noise ratio in order to characterize these photodiodes.

Keywords

Optical communication, simulation, photodetector, semiconductor, *PIN* photodiode, *APD* photodiode, *Optisystem*.

1. INTRODUCTION

Today the new world of telecommunications represents a colossal circle of information where the evolution is incredibly high thanks to the actual communication systems which knew an effective extension of service diversification [1]. The application of optoelectronics in such systems involves optical methods which permit to convey and treat luminous signals using pure optical means; this manner presents numerous advantages comparing with the possibilities offered by electronics, for the simple reason that it is relatively easy to propagate luminous pulses with a very short duration of scale of the picoseconds, this property is attractive to realize more rapid and reliable systems “all-optics” involving radical changes in telecommunications industry [2].

Considered as a fundamental component in the reception module of the optical communication chain, the photodiode (*PIN* “*Photodiode Intrinsic Negative*” or *APD* “*Avalanche PhotoDiode*”) plays a very important role in detecting luminous wave carrying information along the optical fiber, then transforms it into an electrical signal which will be

easily exploited in the extraction of the transmitted information. In this purpose, several studies were being developed aiming to examine the major limitations caused in such optoelectronic component “the photodiode” by processing their physical origins.

In this context, *B. Orsal* and *R. Alabedra* [3] studied the photodiode characteristics by the measurements of $1/f$ noise (*Flicker* noise which belongs to the excess noises). Using the 2nd optical window ($\lambda = 1300$ nm), they found that the *Flicker* noise is due to the dark current for two generations of photodiodes of different contact technology; this noise analysis showed that the quality of contacts of the 2nd generation devices is better than the first generation. On the other hand, they analyzed [4] too the ionization by impact process in the *APD* photodiode (avalanche effect process) which has a relationship with noise of the multiplication introduced by holes which are the less ionizing carriers.

Otherwise, *U. Ibrahim* [5] evaluated and compared between the *PIN* photodiode and *APD* photodiode in terms of performance in an optical multiplexed transmission link using *WDM* “*Wavelength Division Multiplexing*” technique, he demonstrated that the capacity to transport data along 15 km can attain up to 5 Gbit/s per user if the *APD* photodiode is used in the receiver module.

P. Gopal [6] compared also between *PIN* and *APD* performances in atmospheric turbulences of a satellite system for different modulation positions, he found that the performance of the *APD* photodiode is better than that of the *PIN* photodiode.

In the view of this importance, we are interested in the analysis of fundamental photodetection noises [7] generated in *PIN* and *APD* photodiodes consulted in an optical high debit transmission chain simulated by the optical software “*Optisystem*” [8].

This chain is considered as an environment for tests and observations where the objective is to evaluate *PIN* and *APD* photodiodes performances in terms of their functional parameters to practically compare between these two different photodiode types used as detectors for the actual optical communication systems. In fact, this work presents mainly:

- First, a description of photodetection principle created in semiconductors by explaining the optical-electrical interaction which is characterized by the “*photons-electrons*” interface.

- Second, a specific view about *PIN* and *APD* photodiodes which are responsible for photons-electrons conversion by presenting their structure, principle and functional parameters.

- Third, a deep study about fundamental photodetection noises summarized in *Johnson* noise, *Shot* noise and noise of multiplication by exposing the different physical sources that involve such phenomena in *PIN* and *APD* photodiodes.

- Fourth, interpretation and explanation of simulation results implemented by the *Optisystem* software through the optical high debit communication system chosen to analyze the photodetection noises and evaluate the *PIN* and *APD* photodiodes performances in function of signal-to-noise ratio (*SNR*) in order to provide new perspectives for the future transmission axes.

2. SEMICONDUCTORS AND PRINCIPLE OF PHOTODETECTION

A semiconductor is a continuous crystalline environment characterized by its energetic band structure [9]: this energetic structure comprises a valence band which corresponds to the binder molecular states and a conduction band which includes the anti-binder molecular states, these two bands are separated by an energy range situated between the top of the valence band and the bottom of the conduction band known by band gap or energy gap.

The photodetection constitutes the core of optical communication systems because in an optical transmission system, the photoelectrical reception interface is the responsible for optimally detecting the luminous signal emitted by the optical source, reflected and guided by the optical fiber walls. This role is ensured by the photodiode which acts as a photon counter and a current generator.

The photodetection in semiconductors [7, 9] operates on the creation of electron-hole pairs under the action of light: when a semiconductor is illuminated by a photon of an energy higher than or equal to its gap ($h\nu \geq E_g$), the absorbed photon carries the electron from the valence band to the conduction band (excited state), where it will be free to move in the crystalline network under the influence of an electric field, and the two generated carriers (*P* and *N*) are collected in a form of a photocurrent I_{ph} (the number of electron-hole pairs is equal to the number of absorbed photons). The photodiode converts directly the incident photons into electrons which can be either ejected or liberated thanks to the photosensitive semiconductor properties that make the photodiode very fast in terms of photoelectrical conversion. This has practically developed the optical transmission systems for the very long distances thanks to the use of *PIN* and *APD* photodiodes. Fig 1.

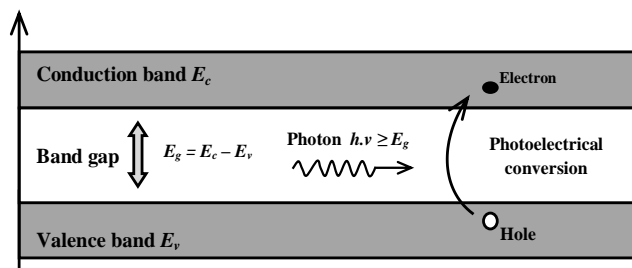


Fig 1: Semiconductor structure and photodetection principle “optical-electrical interaction”

3. PIN AND APD PHOTODIODES

The photodiode is a semiconductor photodetector that possesses two electrodes and a symmetric electrical characteristic [10]. It is used in optical transmissions for its high rapidity and responsivity, indeed, the most important structures disposed to convert the optical flow to an electrical current are:

3.1 PIN photodiode

The *PIN* photodiode [11] is constituted of three zones: the first is *P*-doped in order to create a hole excess, the second presents an intrinsic region “*i*” called absorption zone and the third is *N*-doped in order to create an electron excess, from this the component name is formed. The width of the intrinsic zone is optimized in a way that maximizes *PIN* performances, so, it is preferred that this region be large from which the interest of *PIN* photodiode.

In this context, in order to increase the photodiode responsivity, an anti-reflection dielectric coating is added to the photodetector device by consequence the majority of photons will be absorbed in the intrinsic zone. The photodiode response must be also very fast, it is limited by two factors: junction capacitance and transit time, so, it can be improved by increasing the reverse bias voltage which reduces these two factors [10]. Fig 2.

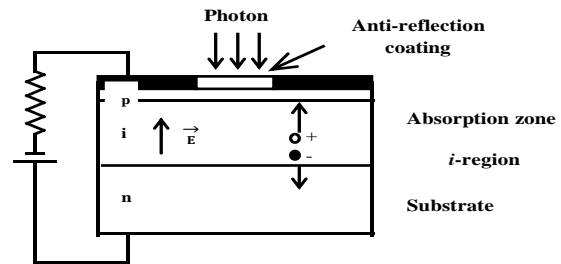


Fig 2: PIN photodiode structure

3.2 APD photodiode

In the *APD* photodiode, the created electrons are multiplied by the avalanche effect (multiplication process) in the intrinsic zone “*i*” in order to obtain an important signal-to-noise ratio, the idea of using this internal multiplication phenomenon is applied so that the incident photon can generate several photoelectrons; this is to increase the electrical signal power which corresponds to the incident optical power. In fact, the *APD* photodiode is just a *PIN* photodiode conceived from a physical ionization by impact process in order to gain a multiplication effect. As a disadvantage, the response time of *APD* photodiode becomes longer due to the additional time occurred by the ionization phenomenon [11, 12]. Fig 3.

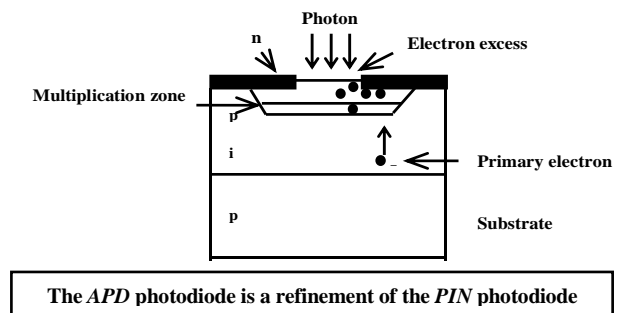


Fig 3: APD photodiode structure

PIN and *APD* photodiodes are characterized by numerous functional parameters, among the main parameters which define such photodetectors there are [7]:

- **Quantum efficiency “ η ”:** It is a very important parameter which has relationship with the probability of creating electron-hole pairs through absorbed photons. It is defined as the fraction of incident photons (absorbed photons) which contribute to the external photocurrent. It can't be higher than 1, and it increases with the thickness of the absorption zone.

$$\eta = \frac{I_{ph}/e}{P_{opt}/h\nu} \quad (1)$$

(I_{ph} : Photocurrent, P_{opt} : Incident optical power, e : Elementary electron charge, h : Planck's constant $6.626 \cdot 10^{-34}$ J.s, ν : Luminous wave speed).

- **Absorption coefficient “ α ”:** One of the most important elements that determine the photodiode performances is the semiconductor absorption which is characterized by its absorption coefficient α which depends on the luminous wavelength λ and of the energy gap E_g . If the photon energy $h\nu$ is higher than the semiconductor gap, there will be absorption, it is necessary then that the luminous wavelength λ be less than the cut-off wavelength λ_c [13, 14].

$$dP_{abs}(z) = P_{opt} e^{-\alpha(\lambda)z} dz \quad (2)$$

(P_{abs} : Absorbed power, α : Absorption coefficient, λ : Luminous wavelength, z : Semiconductor depth)

- **Responsivity “ R ”:** It is a very important parameter which is defined by the ratio of the photocurrent I_{ph} on the received energy flux of photons (incident optical power) [13, 14].

$$R(\lambda) = \frac{I_{ph}}{P_{opt}} = \eta \frac{\lambda}{1,24} A/W \quad (3)$$

- **Noise Equivalent Power “ NEP ”:** The noise equivalent power is defined by the incident optical power for which the signal-to-noise ratio is equal to 1 ($SNR = 1$), in other words, the photocurrent I_{ph} is equal to the noise current I_N . This parameter is inversely related with the photodiode directivity.

$$NEP = \frac{I_N}{\frac{I_{ph}}{P_{opt}}} = D^{-1} \quad (4)$$

- **Gain “ G ”:** It depends on the number of detected carriers in relation with photogenerated electron-hole pairs.

$$G = \frac{I_{ph}}{e \cdot \eta \cdot \frac{\lambda}{h.c} \cdot P_{opt}} \quad (5)$$

- **Flux “ Φ ”:** It is defined by the number of photons reaching into the active photodiode surface at a given wavelength per a unit of time.

$$\Phi = P_{opt} \cdot \frac{\lambda}{h.c} \quad (6)$$

4. NOISE OF PHOTODETECTION

At the reception, because of the corpuscular nature of electromagnetic luminous rays, absorption process of photons will be discrete and carries creation will be contingent in terms of time providing a random process, by consequence, the detection of incident optical rays becomes very inhibited because of photodetection noise. In this sense, the noise of photodetection [7] is generated in form of random intrinsic fluctuations which perturb the useful signal carrying information that makes them extraction very difficult at the photodetector output. In fact, the measurement of the photodetection noise in the presence of luminous signals is based on the evaluation of signal-to-noise that expresses the noise associated to the photocurrent generated through the conversion of photons into electron-hole pairs.

$$SNR = \left(\frac{S}{B} \right)_{dB} = 10 \log \frac{P_{opt}}{NEP} \quad (7)$$

There are five types of photodetection noise classified into two categories:

- **Fundamentals noises:** First, thermal noise “*Johnson noise*” which is due to the random collisions of carriers with crystal atoms which permanently vibrate that introduces a thermal agitation. Second, photonic or quantum noise “*Shot noise*” which is due to the collisions of incident photons. Third, noise of multiplication “*Multiplicative noise*” which is essentially a quantum noise amplified by the avalanche effect (multiplication process).

- **Excess noises:** First, noise with $1/f$ called also “*Flicker noise*” which is associated to the presence of potential barriers in the electrical contacts. Second, noise of generation-recombination “*G.R. noise*” which is due to the random generation and recombination of carries because of gap traps.

In this context, the objective of this research is especially focused on the study of fundamental noises caused in the photodiode:

4.1 Thermal noise “*Johnson noise*”

This noise is the result of random carriers fluctuations caused in the photodiode resistance providing a thermal agitation, because the *Brownien* electron movement produces a random voltage at the resistance extremities.

$$S_i(f) = \frac{4 \cdot k_B \cdot T}{R} \quad (8)$$

(k_B : Boltzmann's constant $1.381 \cdot 10^{-23}$ J/K, T : Temperature in Kelvin, R : Photodiode resistance)

The spectral density of current is independent of the frequency because it is the case of a white noise. In the

photodiode, shunt resistance and series resistance are the two elements that represent the source of the thermal noise.

4.2 Quantum noise “Shot noise”

Shot noise or called *Schottky* noise depends on the corpuscular nature of electrons and their random emission and collection in terms of time:

$$S_i(f) = 2.e.I \quad (9)$$

It is also a white noise where the spectral density of current is uniform. It is due to the discrete fluctuations caused by the current containing electrons which are randomly distributed. In reality, this noise is caused by the dark current of the photodiode I_D and the photocurrent I_{ph} which are generated from the detectable luminous signal (absorbed photons).

$$S_i(f) = 2.e.[I_D + I_{ph}] \quad (10)$$

In the actual components, the noise is essentially due to the photocurrent I_{ph} :

$$I_D \ll I_{ph} \Rightarrow S_i(f) = 2.e.I_{ph} \quad (11)$$

4.3 Noise of multiplication “Multiplicative noise”

The noise of multiplication is specifically generated in APD photodiode possessing an internal gain based on the multiplication method by impact of carriers in depletion region (junction region or space charge region). It is just a quantum noise caused in PIN photodiode which will be amplified in APD photodiode by the avalanche effect:

$$S_i(f) = 2.e.I_{pho} \cdot \bar{M}^2 \quad (12)$$

(\bar{M} : Average value of multiplication factor of APD photodiode, I_{pho} : Primary photocurrent)

5. SIMULATION RESULTS

The simulation part was executed using the *Optisystem* software which constitutes an interactive environment that permits to simulate, observe and analyze the transmitted signal at the level of all modules of the optical high debit transmission chain [8] considered as a model for the study in a form of schematic blocks as displays the figure 4, basing on the following parameters:

- Emitted power $P_e = 50$ mW, binary debit $D = 10$ Gbit/s, laser diode wavelength $\lambda = 1552.52$ nm corresponding to the 3rd optical window, mono-mode fiber length $L_{Fib} = 50$ Km, photodiode sensitivity $S = 0.8$ A/W, photodiode dark current $I_D = 5$ nA.

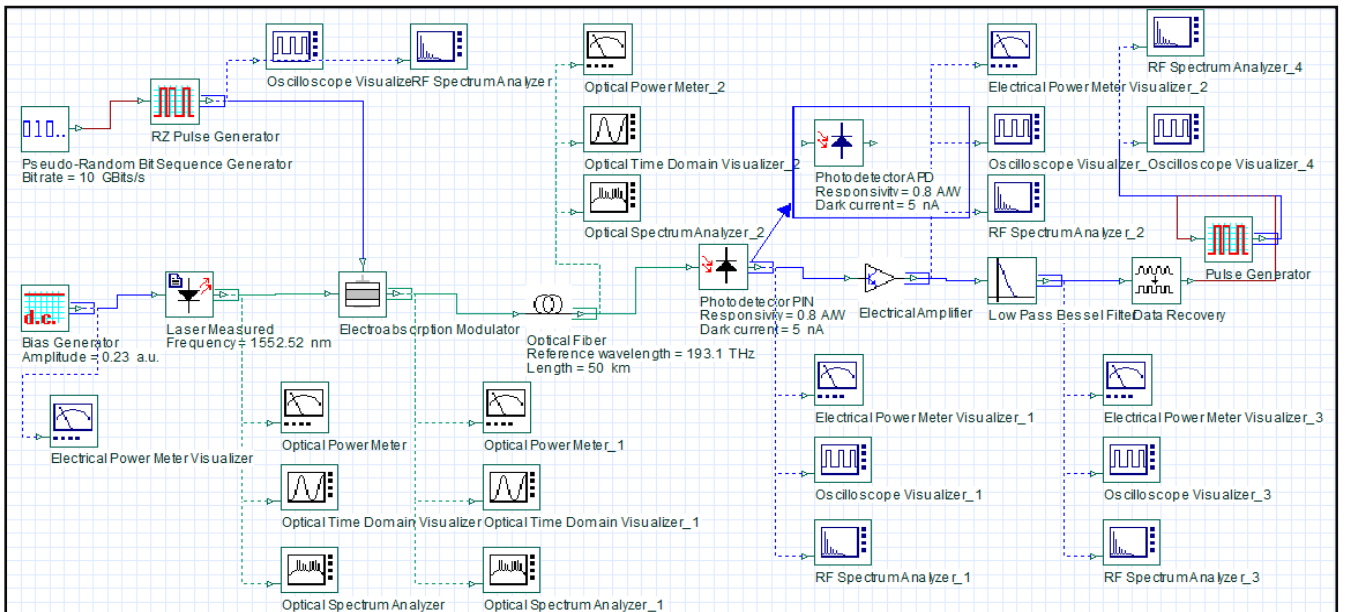


Fig 4: Model of simulation realized by Optisystem: Optical high debit transmission chain “ $P_e = 50$ mW, $D = 10$ Gbit/s, $\lambda = 1552.52$ nm, $L_{Fib} = 50$ Km, $S = 0.8$ A/W, $I_D = 5$ nA, $G_{APD} = 3dB$ ”

The aim of this study is to examine the noise of photodetection caused in PIN and APD photodiodes used as detectors in the transmission chain basing on a temporal visualization (in terms of time using an optical time domain visualizer) or a spectral visualization (in terms of frequency using an optical spectrum analyzer).

In specific, it is distinguished that the photodiode constitutes the seat of the noise which is observed additive to the useful signal; this noise has a random character manifested by parasitic fluctuations that distort the electrical pulses containing

information as shows the figure 5. It is the noise of photodetection [15] whose the sources are internal generated in the photodiode core; this noise has a low power but it equivocally influences the received signal consequently the transmitted information.

The objective of this study extends to define the dominant noise contribution in both PIN and APD photodiodes considered as detectors for the studied transmission chain in order to compare and evaluate their performances, in fact, this simulation consists of two major steps:

5.1 Analysis of photodetection noise in PIN and APD photodiode

This first step consists to simulate the optical chain using respectively the *PIN* photodiode and the *APD* photodiode in order to visualize the shape of thermal and quantum noises caused in these photodiodes using a power spectrum analyzer.

As displays the figure 6 (a, b), the thermal noise power caused in the *PIN* photodiode is higher than it in the *APD* photodiode.

On the contrary, the quantum noise power caused in the *APD* photodiode is denser than it in the *PIN* photodiode as displays the figure 7 (a, b).

The dominant noise in the *PIN* photodiode is that of thermal noise derived from the various circuits contained in the photodiode such as the resistance and the preamplifier which cause thermal agitation of electrons which consequently influences on the photodiode responsivity, however the *APD* photodiode resists against the thermal noise thanks to its internal

gain produced by the avalanche process, but this process causes a rising of quantum noise that reduces the photodiode detection speed, this is due to the noise of multiplication generated particularly in *APD* photodiode by impact of carriers in the depletion region. In this view, the optical system was simulated again using an *APD* photodiode by varying its internal gain from 2 to 15 dB. The figure 8 presents both the detected signal and the noise of photodetection measured by an electrical power measurer, and the signal-to-noise ratio calculated from the two previous values in function of the avalanche factor.

It is observed that more the internal *APD* gain increases more the detected signal and the photodetection noise increase, this involves a degradation of signal-to-noise ratio caused by the quantum noise which is amplified more rapidly than the detected signal. This phenomenon is due to the noise of multiplication also known as avalanche noise because it is only generated in *APD* photodiode [4].

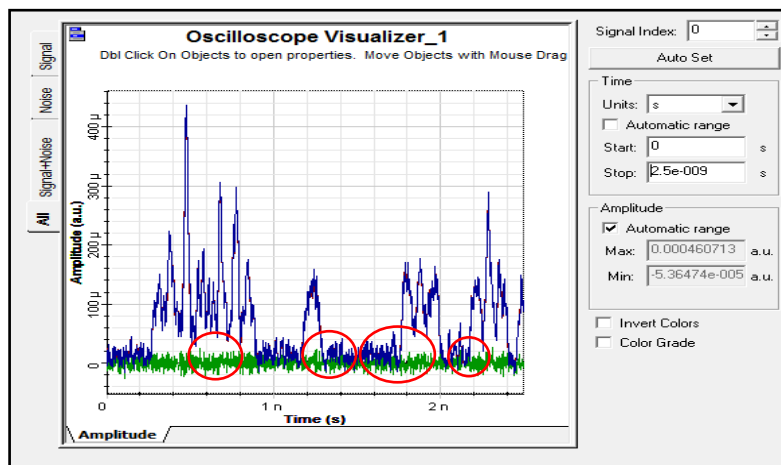
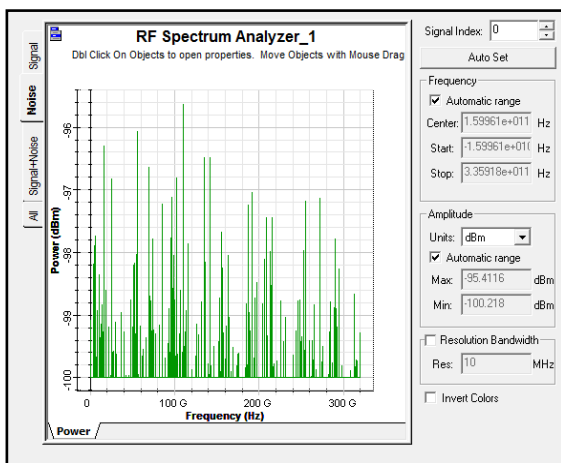
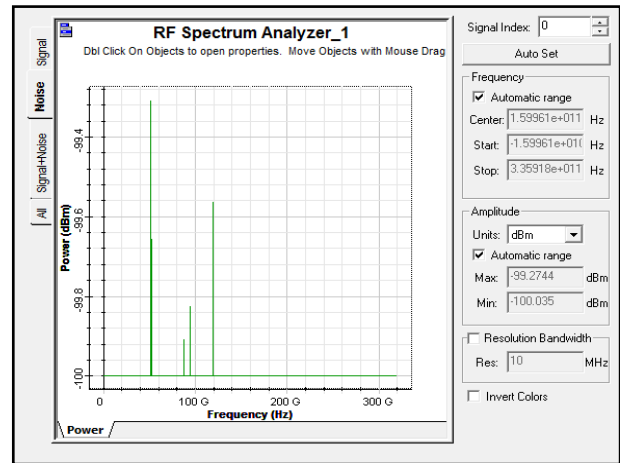


Fig 5: *PIN* photodiode output (detected signal in blue) – Noise of photodetection (noise signal in green) – Fluctuations (distorted signal surrounded by the red circles)

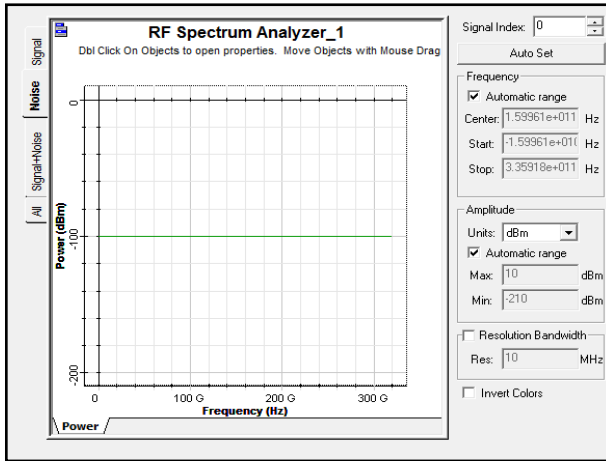


(a) *PIN* photodiode

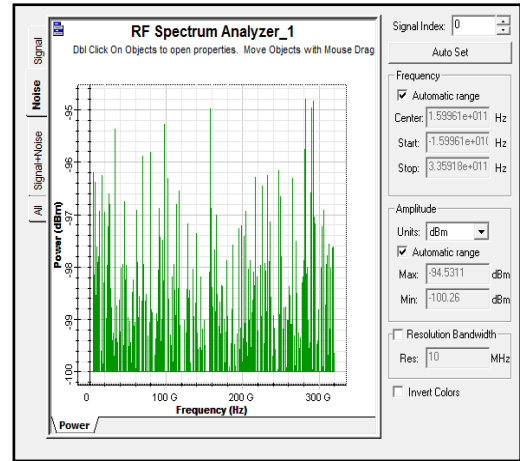


(b) *APD* photodiode

Fig 6: Spectral visualization of thermal noise “Johnson noise” caused in *PIN* and *APD* photodiodes



(a) PIN photodiode



(b) APD photodiode

Fig 7: Spectral visualization of quantum noise “Shot noise” caused in PIN and APD photodiodes

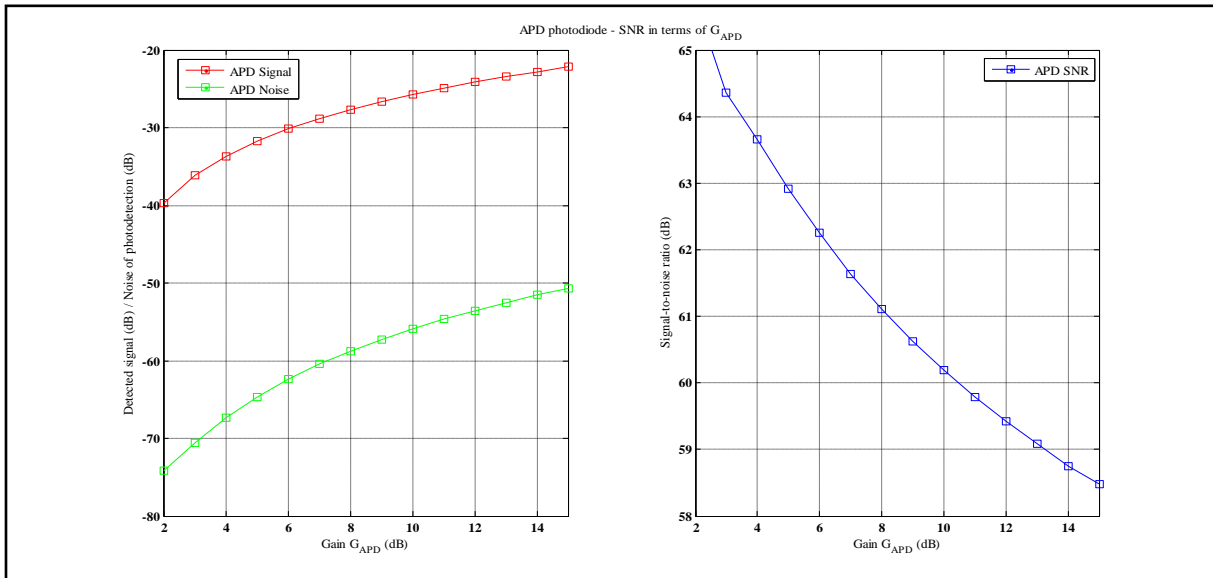


Fig 8: Signal-to-noise ratio in terms of internal gain of the avalanche photodiode APD – $SNR = f(G_{APD})$

5.2 Evaluation of performances of PIN and APD photodiodes

The objective of this second step is to study the performance of PIN and APD photodiodes according to the following functional parameters: emitted power P_e , debit D , responsivity R and dark current I_D by measurement of the detected signal issued at the photodiode output and the photodetection noise generated in the photodiode core by using the electrical power measurer and then calculate the signal-to-noise ratio in decibel [16].

In this step a binary debit of 2.5 Gbit/s was contributed with using for the first time a PIN photodiode and for the second time an APD photodiode by varying in every time only one of the previous parameters for each photodiode.

The results are represented in the following curves shown respectively in the figures 9, 10, 11, 12:

- **In terms of the emitted power - $S/B = f(P_e)$:** In the figure 9, it is clearly observed that more the input power “ P_e ” increases more the detected signal is better, however the noise is relatively constant and seems lower than the received signal. This expresses that the signal-to-noise ratio rises with the

emitted power. Besides, the comparison between the PIN and APD photodiodes justifies that the APD photodiode is more advantageous than the PIN photodiode for the low powers, by contrast the PIN photodiode remains always more efficient for the high powers upper than 0.7 W [17].

Subsequently, the emitted power will be equal to 50 mW, a value which belongs to the low powers range and justifies that the APD performances will be better than the PIN performances in terms of signal-to-noise ratio for the following sub-parts.

- **In terms of debit - $S/B = f(D)$:** The figure 10 shows that the input binary debit “ D ” varies from 100 Mbit/s to 55 Gbit/s, it notes that more the input debit increases more the photodetection noise increases but the detected signal is relatively attenuated. In fact, raising the data debit reduces the signal-to-noise ratio that decreases therefore the transmission performances. Finally, if we need to keep a good performance for the studied system, it should not exceed a debit of 25 Gbit/s, because above this value, the signal-to-noise ratio fall and becomes not acceptable.

- **In terms of responsivity - $S/B = f(R)$:** In the figure 11, the responsivity “ R ” varies from 0.1 to 1 A/W, it is directly

observed that more the responsivity increases more the detected signal increases, however the photodetection noise takes very low values that makes the signal-to-noise ratio better consequently the transmission performances. This demonstrates that the signal-to-noise ratio is in a good agreement with the responsivity of the photodiode, *PIN* or *APD* and becomes adequate above a value of 0.3 A/W.

- **In terms of dark current - $S/B = f(I_D)$:** The figure 12 shows that the dark current " I_D " varies from 0 to 10000 nA. It is simply observed that more the dark current increases and takes higher values, the detected signal, the noise of photodetection and the signal-to-noise ratio keep always constant values. This proves that the dark current is negligible for the photoconductor mode [7], its value is very low to be considered in the generation of the noise of photodetection that confirms what was already mentioned in the section (4, b).

Finally, the *APD* photodiode retains always the better performance comparing with the *PIN* photodiode for the reason that the manipulation used particularly a low input power equal to 50 mW.

6. CONCLUSION

This work has mainly aimed to analyze and describe the fundamental photodetection noises caused in *PIN* and *APD* photodiodes; this permitted to study the physical sources generating the noise of photodetection which displays internal and random, and noticeably affects the photodiode signal. This research has also aimed to characterize these photodetectors by examining their performances in terms of different functional parameters by measurement of the signal-to-noise ratio which constitutes a very important factor to determine the quality of the optical detection for the optical high debit communication chain chosen as a model of simulation and study; this led to conclude that these parameters constitute always major factors in evaluating *PIN* and *APD* performances, making their choice more judicious for high transmission quality in the various optical communication systems having an important role in the new telecommunications world because today we can't speak about telecommunications without mention the optical communication systems characterized by a grand demand for their exceptional quality offering high debit, long distance of propagation and strong immunity against perturbations.

Finally, it is very interesting to extend this research by studying the other photodetection noises which are the excess noises

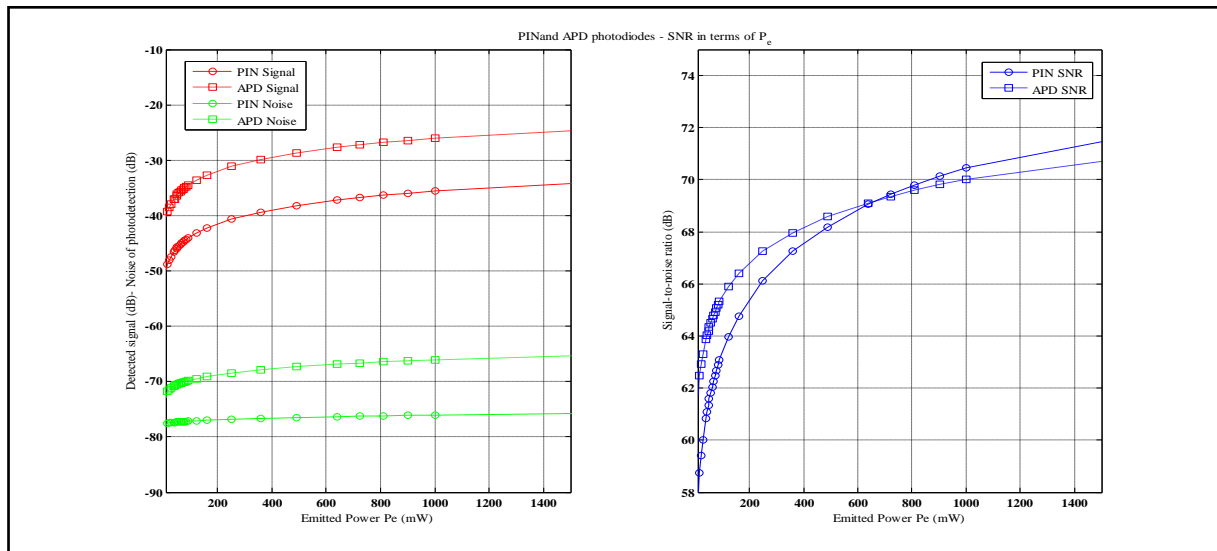


Fig 9: Signal-to-noise ratio in terms of emitted power for *PIN* and *APD* photodiodes – $SNR = f(P_e)$

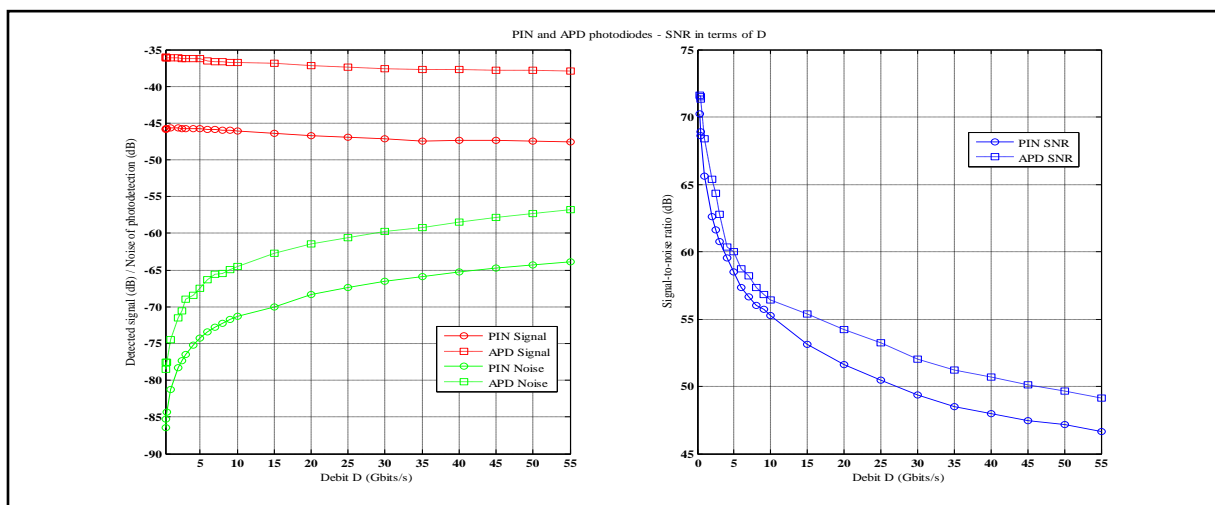


Fig 10: Signal-to-noise ratio in terms of binary debit for *PIN* and *APD* photodiodes - $SNR = f(D)$

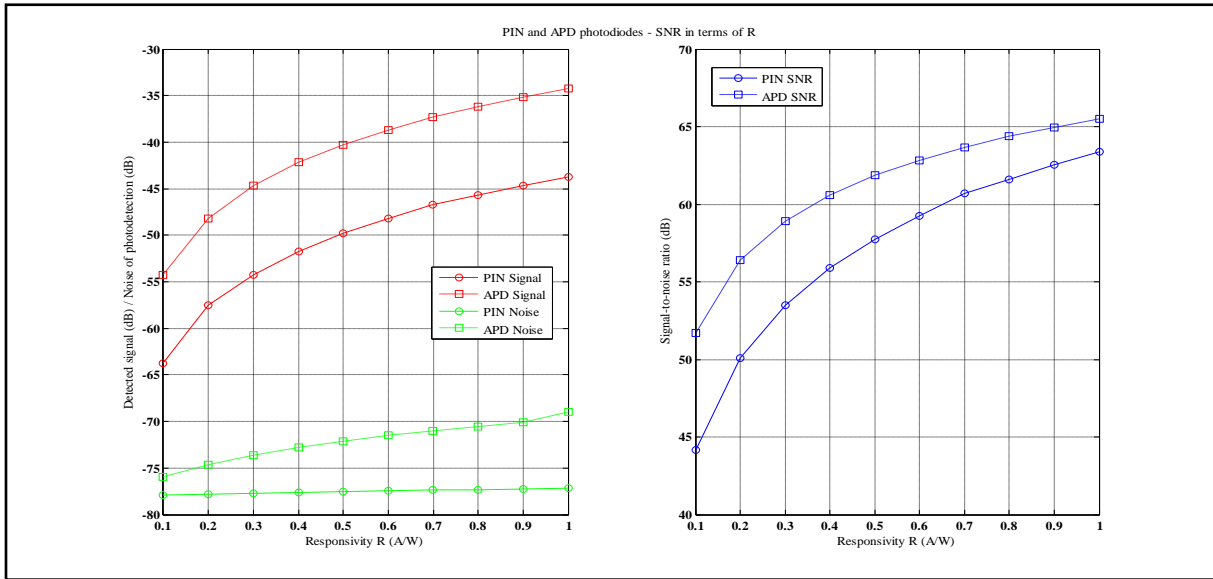


Fig 11: Signal-to-noise ratio in terms of responsivity for *PIN* and *APD* photodiodes – $SNR = f(R)$

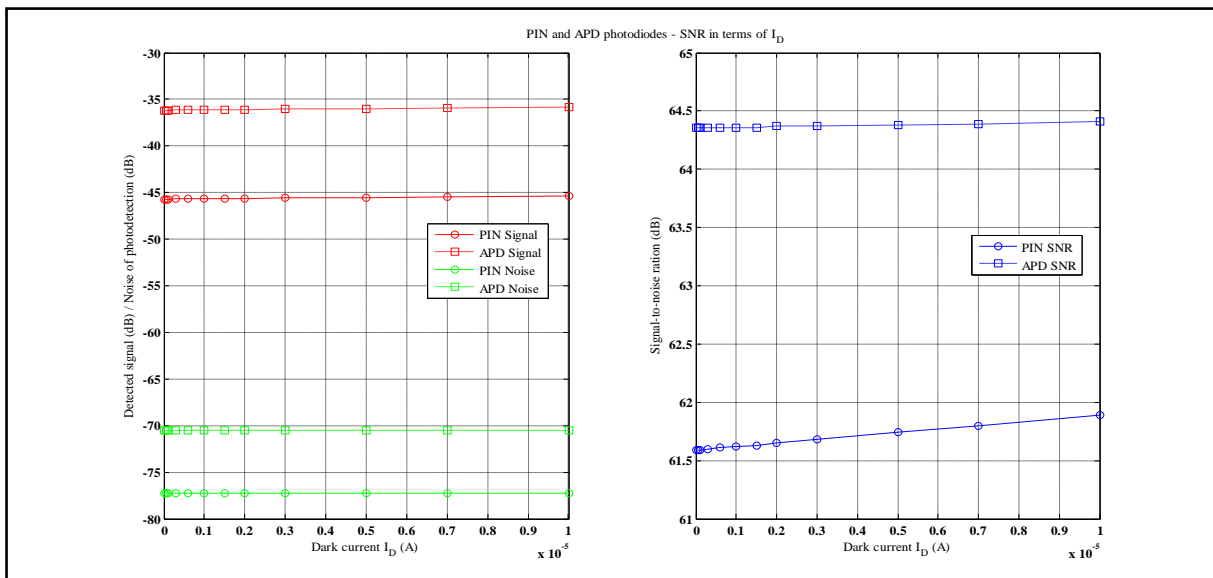


Fig 12: Signal-to-noise ratio in terms of dark current for *PIN* and *APD* photodiodes – $SNR = f(I_D)$

summarized in “Flicker noise” and “generation-recombination noise”, as it is very important to model all types of photodetection noise in order to get a good performance of photodetector that establishes a topical subject making the optical communication systems more efficient.

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